

Reactions produced by μ mesons in hydrogen

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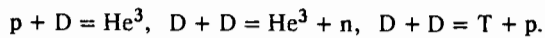
(Submitted 8 February 1957)

Zh. Eksp. Teor. Fiz. **32**, 947-949 (1957) [*Sov. Phys. JETP* **5**, 775-777 (1957)]. Also *S1*, pp. 7-10

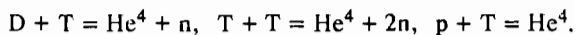
Usp. Fiz. Nauk **161**, 43-46 (May 1991)

At present there is evidence that at Berkeley¹ in a bubble chamber filled with liquid hydrogen having a varying deuterium content, there was observed a nuclear reaction catalyzed by μ mesons. The possibility of such a reaction was first pointed out by Frank² in connection with the analysis of π - μ disintegrations in emulsions. This process was investigated in liquid deuterium by both of the authors of this article, independently of one another.^{3,4}

The presence of a μ meson changes the form of the potential barrier, which previously prevented nuclear reactions among slow proton and deuteron nuclei, increasing sharply the penetrability of the barrier and making possible the reactions



In the presence of tritium there are also possible the reactions



The reaction $p + p = D + e^+ + \nu$, catalyzed by mesons, is practically impossible, since in addition to the barrier there is also the factor of a low probability for the beta process.

It has been predicted⁴ that the probability of the reaction in flight is low, the production of mesomolecules practically always leads to nuclear reactions, the rate of the process is determined by the production of mesomolecules, and the probability of mesomolecule formation during the lifetime of a meson may amount to several hundredths or even tenths, depending on the arrangement of the mesomolecule levels.

The experimental data of Alvarez¹ show that in natural hydrogen (deuterium content ratio 1:7000), the reaction $p + d = \text{He}^3$ occurs on the average once for each 150 mesons. If the deuterium ratio is 1:300, the reaction occurs once per 40 mesons, and if the ratio is 1:20, the reaction occurs once per 33 mesons. Furthermore, the energy of the resulting He^3 (5.4 MeV) is carried away by the μ meson, so that monochromatic μ mesons are observed while the reaction is taking place. The relatively high probability found for the reaction in the natural mixture is explained¹ by the transfer of the meson from the proton to the deuteron (charge exchange): $p\mu + d = p + d\mu$. Because of the difference in reduced mass, the energy of the $D\mu$ bond (2655 eV) is greater by $\Delta E = 135$ eV than the energy of the $p\mu$ bond. Therefore the charge-exchange process appears to be irreversible under the experimental conditions.

We shall give a rough estimate of the probability of the transition. If ΔE is equal to zero, the cross section should be of the order of πa^2 , where a is the radius of the Bohr orbit of the mesoatom, 2.5×10^{-11} cm.

Indeed, if the masses of the two nuclei are equal, with $\Delta E = 0$, the states of the systems Σ_{+g} and Σ_{-u} appear to be proper, and the cross section for charge exchange can be expressed by the scattering lengths a_g and a_u of these states in a continuous spectrum: $\sigma = \pi(a_g - a_u)^2$. When $\Delta E \neq 0$, but is still small with respect to the molecular dissociation energy, then

$$\sigma = \pi(a_g - a_u)^2 v_i / v_f,$$

where v_i is the velocity before collision and v_f is the velocity after collision.

In actual fact, ΔE is of the same order as the dissociation energy, so that the formula is corrected at least in order of magnitude. If v_i is small, $\sigma \sim 1/v_i$. It follows that in order of magnitude,

$$\sigma \approx \pi a^2 v_* / v,$$

where v_* is the characteristic velocity corresponding to 1.35 eV. Using the masses of the proton and deuteron, $v_* = 2 \times 10^7$ cm/sec. Such an estimate yields qualitative agreement with the observed facts. Calculation shows that saturation should be reached at deuterium concentrations of 1 in 300 to 1 in 20. In the natural mixture of hydrogen, the probability of $D\mu$ production and consequently the probability of the reaction should be threefold less than in enriched mixtures; experimentally it is found to be four to five times less.

Let us examine the reaction in the $pD\mu$ molecule. The observed high probability a of the giving up of meson energy does not agree with the hypothesis that the reaction proceeds like an electric dipole transition $E1$, since a becomes 2×10^{-3} for the meson under this hypothesis. Therefore, to estimate the probability we cannot use the experimental cross section for $p + D = \text{He}^3 + \gamma$ (as was done before⁴), since under the conditions of the measurement, it is precisely the cross section of the $E1$ process which is observed.⁵⁻⁶

In the case of zero orbital momentum, the system $p + D$ can be found in either of the states $+3/2$ or $+1/2$. The transition to He^3 ($+1/2$ state) is possible in the first case as $M1$ and $E2$, and in the second case as $M1$ and $E0$.¹¹

The conversion coefficients have the following values: for $M1$, $a = 10^{-4}$; for $E2$, $a = 0.1$; for $E0$, only the giving up of meson energy is possible (the probability of pair production e^+ and e^- at the expense of $E0$ is 10^{-3} of the probability of the giving up of meson energy in the case of $p + D$, but is of the order of unity in the case of $p + \text{H}^3 = \text{He}^4$).

Calculations concerning barrier penetration for the $pD\mu$ molecule under adiabatic investigation of the motion of the proton and deuteron yield $\psi^2(0) = 6 \times 10^{-27}$ cm⁻³.

For the mirror reaction $n + D$, it is assumed that for the



FIG. 1. Ya. B. Zel'dovich and A. D. Sakharov as guests of D. A. Frank-Kamenetskiĭ (on the right), 1950's, Arzamas-16(Sarov).

purposes of computation⁸ that the process proceeds from the state $3/2$ as the result of M 1. Experimentally for thermal neutrons,⁹ $\sigma = 5.7 \times 10^{-28}$ cm² with $v = 2200$ m/sec and $\sigma v = 1.3 \times 10^{-22}$ cm³/sec.

Hence for mesomolecules the probability of the reaction ($\tau = 2.15 \times 10^{-6}$ being the meson lifetime) is

$$w = \frac{\sigma v \psi^2(0)}{(1/\tau) + \sigma v \psi^2(0)} = 0.6.$$

During the approach of the proton and deuteron in the spin state $1/2$, the approximate determination of the magnitude of the monopole moment was carried out by examining one charged particle with the wave function $\psi = (1/\sqrt{2\pi\lambda r})e^{-r/\lambda}$ in the final (combined) state and $\psi = \psi(0)(1 - \lambda/r)$ in the initial state [continuous spectrum, $\psi(0)$ being the previously calculated wave function under the barrier, $\psi^2(0) = 6 \times 10^{-27}$ cm⁻³].

The probability of a process with a release of energy, to a meson in a $p + D$, in spin state $1/2$, and with $\lambda = 2.4 \times 10^{-13}$ cm, turned out to be equal to 0.5.²⁾

Thus from the rough estimates made above, it follows that the probability observed by Alvarez for the process which involves release of energy to the meson, and the probability of the process which involves emission of a gamma quantum, can both be close to unity during the meson lifetime.

In a more accurate investigation, not only would it be necessary to take into account the fact that the process is not adiabatic (thus involving terms of the order of the meson-nucleon mass ratio), but it would also be necessary to make a

separate investigation of the nuclear reaction with different values of total molecular spin.

Note Added in Proof (February 9, 1957)

The probabilities for mesomolecule production in the collisions $D\mu + p = Dp\mu$ and $D\mu + D = D_2\mu$ differ not only because of the different positions of the excited vibrational levels of the molecules,⁴ but also because in $pD\mu$ the center of mass does not coincide with the charge center, and thus possesses a dipole moment $\frac{1}{2}ea$. Therefore in the collision of slow $D\mu + p$ there is possible the dipole transition $E1$ in the molecule into the momentum state 1, with energy given up to the electron. In the case of $D\mu + D$, only the $E2$ transition into momentum state 2 competes with the $E0$ transition investigated.⁴

¹⁾ Church and Weneser⁷ have recently drawn attention to the role of $E0$ in the case of internal conversion in the transition $J \rightarrow J \neq 0$.

²⁾ Choosing $\lambda = \hbar/\sqrt{2ME}$, where M is the reduced mass of p and D , and E is the binding energy 5.4 MeV.

¹ L. W. Alvarez *et al.*, Lithographed document, December 1956.

² F. C. Frank, *Nature* 160:525 (1947).

³ A. D. Sakharov, Report of the Physics Institute, Academy of Sciences, USSR (1948).

⁴ Ya. B. Zel'dovich, *Dokl. Akad. Nauk SSSR* 95:454 (1954).

⁵ G. M. Griffiths and I. B. Warren, *Proc. Phys. Soc.* 68:781 (1955).

⁶ D. H. Wilkinson, *Phil. Mag.* 43:659 (1952).

⁷ E. L. Church and J. Weneser, *Phys. Rev.* 103:1035 (1956).

⁸ N. Austern, *Phys. Rev.* 83: 672 (1951); 85:147 (1952).

⁹ Kaplan, Ringo, and Wilzbach, *Phys. Rev.* 87:785 (1952).

¹⁰ E. E. Salpeter, *Phys. Rev.* 88:547 (1952).

Translated by D. A. Kellogg