

EXPERIMENTAL DATA ON  $\mu \rightarrow e$  DECAY

I. I. GUREVICH and B. A. NIKOL'SKIĬ

I. V. Kurchatov Institute of Atomic Energy

Usp. Fiz. Nauk 95, 476-479 (July, 1968)

1. THE mass, magnetic moment, and the lifetime of the muon are respectively equal to

$$m = 105.659 \pm 0.002 \text{ MeV},$$

$$\mu = (1.001164 \pm 0.000003) \frac{e\hbar}{2m_\mu c},$$

$$\tau = (2.199 \pm 0.001) \cdot 10^{-6} \text{ sec}.$$

2. In the decay of a fully polarized muon, the spectrum of the electrons is given by [1-3]

$$dN(x, \vartheta) = \frac{1}{2} \left\{ \frac{1+f(x)}{1+4\eta \frac{m_e}{m_\mu}} \left[ 42 - 12x + \rho \left( \frac{32}{3}x - 8 \right) + 12 \frac{m_e}{m_\mu} \cdot \frac{1-x}{x} \eta \right] \right. \\ \left. \pm \xi \cos \vartheta \left[ 4 - 4x + \delta \left( \frac{32}{3}x - 8 \right) + g(x) \right] \right\} x^2 dx d(\cos \vartheta).$$

The signs “+” and “-” pertain to  $\mu^+$  and  $\mu^-$  mesons, respectively;  $x$  is the electron momentum measured in units of the maximum value of this quantity;  $\rho$ ,  $\delta$ ,  $\xi$ , and  $\eta$  are parameters which are bilinear combinations of the interaction constant (see Sec. 3);  $f(x)$  and  $g(x)$  describe the corrections to the spectrum for the radiative effects. The experimental values of the parameters  $\rho$ ,  $\delta$ ,  $\xi$ , and  $\eta$ , and also of the degree of polarization of the electrons  $h$  are listed in the table.

3. Even an absolutely accurate knowledge of the six parameters to be determined from experiment ( $\tau$ ,  $h$ ,  $\rho$ ,  $\xi$ ,  $\eta$ , and  $\delta$ ) will not answer the question of the form of the interaction in the  $\mu \rightarrow e$  decay. The most general form of the interaction Hamiltonian of the  $\mu \rightarrow e$  decay is

$$\mathcal{H} = \sum_{k=1}^5 (\bar{e}\Gamma_k\mu) (\bar{\nu}\Gamma_k(g_k + g_k\gamma_5)\nu) + \text{h.c.}, \quad (2)$$

where  $\Gamma_k$  with  $k = 1, \dots, 5$  denotes the operators 1,

$\gamma_\alpha$ ,  $\tau_{\alpha\rho}$ ,  $i\gamma_\alpha\gamma_5$ , and  $\gamma_5$ .

It follows from the Hamiltonian (2) that the theory is determined by ten complex constants  $g_k$  and  $g'_k$ . Eliminating the insignificant common phase factor, we obtain 19 constants to be determined. The six experimental parameters determined in the study of the electron spectrum of the  $\mu \rightarrow e$  decay yield only six relations between the 19 interaction constants.

In the case when  $h \neq 1$ , it is possible to determine experimentally the parameters  $\rho(h)$  and  $\delta(h)$  as functions of  $h$ . Even in this case, however, the constants  $g_k$  and  $g'_k$  cannot be determined uniquely. [15]

A unique determination of the constants  $g_k$  and  $g'_k$  from experiments on  $\mu \rightarrow e$  decay is possible only if the decay neutrino is registered besides the electron. [16] Such experiments, however, are not realistic at present, although they are highly desirable.

4. Experiments on the study of the electron spectrum of the  $\mu \rightarrow e$  decay can be used for a unique determination of the constants of a theory that is less general than the theory determined by the Hamiltonian (2). Let us consider a theory in which a two-component neutrino is assumed. Such a theory leads [17] to experimental parameters  $\rho = \delta \equiv 3/4$  and  $\xi = -h$ . The electron spectrum (1) is then determined by only two parameters,  $\xi$  and  $\eta$ . All that remain in the theory are two (complex) constants, [17] since in the case of a two-component neutrino  $g_k = g'_k = 0$  for the scalar, pseudoscalar, and tensor interactions. For the vector and axial interactions we have for the constants in the Hamiltonian (2)  $g_V = -g'_V$  and  $g_A = -g'_A$ . The equalities  $\eta = 0$  and  $\xi = 1$  correspond to the V-A interaction. The two complex constants  $g_V$  and  $g_A$  are uniquely determined by

Experimental values of the parameters of the electron spectrum of the  $\mu \rightarrow e$  decay [7]

Parameter	Value	Reference	Parameter	Value	Reference
$\rho$	$0.745 \pm 0.025$	4	$h$	$1.05 \pm 0.3$	11
$\rho$	$0.750 \pm 0.003$	5	$h$	$0.94 \pm 0.38$	12
$\rho$	$0.760 \pm 0.009$	6	$h$	$1.04 \pm 0.18$	13
$\rho$	$0.762 \pm 0.008$	7	$h$	$-0.89 \pm 0.28$	14
$\rho^*)$	$0.750 \pm 0.006$	Mean value	$h$	$(-1.00 \pm 0.13)$	Mean value
$\delta$	$0.78 \pm 0.05$	4	$\eta^{**})$	$-2.0 \pm 0.9$	4
$\delta$	$0.782 \pm 0.031$	8	$\eta$	$0.05 \pm 0.5$	5
$\delta$	$0.752 \pm 0.009$	7	$\eta$	$-0.7 \pm 0.6$	6
$\delta$	$0.754 \pm 0.0085$	Mean value	$\eta$	$-0.7 \pm 0.5$	7
$\xi$	$0.97 \pm 0.05$	9	$\eta$	$-0.8 \pm 0.4$	Mean value
$\xi$	$0.94 \pm 0.07$	4			
$\xi$	$0.975 \pm 0.015$	10			
$\xi$	$0.973 \pm 0.014$	Mean value			

\*The error is increased to allow for the fact that the parameter  $\eta$  may differ from zero [7].

\*\*To obtain this value of  $\eta$ , both parameters of the isotropic part of the spectrum  $\rho$  and  $\eta$  were varied. In all other cases the parameter  $\eta$  was calculated under the assumption that  $\rho \equiv 3/4$ .

three experimental parameters,  $\xi$ ,  $\eta$ , and  $\tau$ , since the phase factor common to the two constants is immaterial. The parameters  $\xi$  and  $\eta$  are expressed in terms of  $g_V$  and  $g_A$  as follows:<sup>[17]</sup>

$$\xi = \frac{2 \operatorname{Re}(g_V g_A^*)}{|g_A|^2 + |g_V|^2}, \quad \eta = \frac{|g_A|^2 - |g_V|^2}{|g_A|^2 + |g_V|^2}. \quad (3)$$

The parameter  $\eta$  has so far been determined with poor accuracy, since this parameter enters in the expression for the electron spectrum (1) with a small factor  $m_e/m_\mu$  (see the table). To determine the ratio  $\varepsilon = g_A/g_V = |\varepsilon| e^{i\theta}$  in the two-component neutrino theory it is possible for the time being to use only the first equation of (3):

$$\xi = \frac{2 \operatorname{Re} \varepsilon}{1 - |\varepsilon|^2} = \frac{2 |\varepsilon| \cos \theta}{1 - |\varepsilon|^2}. \quad (4)$$

Assuming the experimental value  $\xi_{\min} = 0.975 - 0.015 = 0.96$ ,<sup>[10]</sup> we get for  $|\varepsilon|$  the estimate  $0.75 \leq |\varepsilon| \leq 1.34$ . We can obtain from (4) also the estimate  $\theta \leq 16^\circ$ . Formula (4) is very insensitive to the values of  $|\varepsilon|$  and  $\theta$ , in spite of the high accuracy with which  $\xi$  is determined. It follows from (4) that the maximum deviation of  $|\varepsilon|$  from unity is obtained in the T-invariant theory ( $\theta = 0$ ).

<sup>1</sup>C. Bouchiat, L. Michel, Phys. Rev. 106, 170 (1957).

<sup>2</sup>T. Kinoshita, A. Sirlin, Phys. Rev. 113, 1652 (1959).

<sup>3</sup>S. M. Berman, Phys. Rev. 112, 267 (1958).

<sup>4</sup>R. Plano, Phys. Rev. 119, 1400 (1960).

<sup>5</sup>M. Bardon, P. Norton, J. Peoples, A. M. Sachs, and J. Lee-Franzini, Phys. Rev. Lett. 14, 449 (1965); J. Peoples, Nevis Report, 147 (1966).

<sup>6</sup>B. A. Sherwood, Phys. Rev. 156, 1475 (1967).

<sup>7</sup>D. Fryberger, Preprint 67-51, Chicago, USA (1967).

<sup>8</sup>H. Kruger, University of California, Report UCRL-9322 (1961).

<sup>9</sup>M. Bardon, D. Berley, and L. M. Lederman, Phys. Rev. Lett. 2, 561 (1959).

<sup>10</sup>V. V. Akhmanov, I. I. Gurevich, Yu. N. Dobretsov, L. A. Makar'ina, A. P. Mishakova, N. A. Nikol'skiĭ, B. V. Sokolov, L. V. Surkova, and V. D. Shestakov, Yad. Fiz. 6, 316 (1967) [Sov. J. Nucl. Phys. 6, 230 (1968)].

<sup>11</sup>A. Buhler, N. Cabbibo, M. Fidecaro, T. Massam, Th. Müller, M. Schneegans, and A. Zichichi, Phys. Lett. 7, 368 (1963).

<sup>12</sup>S. Bloom, L. A. Dich, L. Feuvrais, G. R. Henry, P. C. Macq, and M. Spighel, Phys. Lett. 8, 87 (1964).

<sup>13</sup>J. Duclos, J. Heintze, A. de Rujula, and V. Soergel, Phys. Lett. 9, 62 (1964).

<sup>14</sup>D. M. Schwartz, Phys. Rev. 162, 1306 (1967).

<sup>15</sup>T. Kinoshita, A. Sirlin, Phys. Rev. 108, 844 (1957).

<sup>16</sup>C. Jarlskog, Nucl. Phys. 75, 659 (1966).

<sup>17</sup>T. Kinoshita and A. Sirlin, Phys. Rev. 107, 593 (1957).

Translated by J. G. Adashko