NEW MAGNETIC MATERIALS-GARNET FERRITES

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WITHIN the last year and a half or two years, great interest has developed in the study of a new class of ferrites that have the so-called garnet structure. This interest is caused by the fact that the garnet ferrites display properties favorable for their use in microwave applications. Their main advantage, as compared with ferrites with spinel structure, is an extremely narrow ferromagnetic resonance-absorption line. Crystals of some garnet ferrites give an absorption line-width of a few oersteds,^{1,2} whereas crystals of spinel ferrites have a width of several tens of oersteds.

Apart from the practical aspect, the study of garnet ferrites is of great interest for the theory of magnetism and crystal chemistry, because their structure is quite different from the familiar structure of the spinel ferrites. Knowledge of their ferromagnetic properties gives additional information about the nature of spontaneous magnetization and its relation to the structure of the material.

The ferromagnetic properties of the garnet ferrites were discovered quite recently, in 1956, by the French scientists Bertaut and Forrat³ and Pauthenet,⁴ and independently a little later, in 1957, by the Americans Geller and Gilleo. Garnet ferrites have the formula

$3 \operatorname{Me}_2 \operatorname{O}_3 \cdot 5 \operatorname{Fe}_2 \operatorname{O}_3 (\mathbf{Or} \operatorname{Me}_3 \operatorname{Fe}_2 \operatorname{Fe}_3 \operatorname{O}_{12}),$

where Me is one of the rare-earth elements or yttrium.

The preparation of such a type of compound is carried out in accordance with usual ceramic practice, from iron oxides and oxides of rare-earth elements. Crystals are grown from the melt by slow cooling from 1325°C to about 900°C, with lead oxide as solvent.⁶ The crystal lattice of a garnet ferrite is cubic; it contains a very large number of ions located at different crystallographic sites. Because of the complexity of the distribution of anions and cations within it, it is almost impossible to depict. The Fe^{3+} and Me^{3+} cations occupy spaces between the oxygen ions. These can be fourfold, the 24 d sites (tetrahedral spaces); sixfold, the 16 a sites (octahedral spaces); or eightfold, the 24 c sites. The Fe^{3+} ions are distributed over the d and a sites, the Me³⁺ ions over the c sites.

To explain the magnetic properties of the garnet ferrites, Néel⁷ proposed to regard the lattice of these substances as composed of three sublattices (in contrast to the spinels, in which, of course, two sublattices are considered, with more octahedral than tetrahedral sites). The Fe^{3+} ions are distributed over two of these (the a and the d), with more tetrahedral than octahedral sites. The Me^{3+} ions are distributed over the third sublattice (the c). According to Néel, there is a strong negative interaction between the first two sublattices, a and d; as a result, there occurs a resultant spontaneous magnetization J (on account of the excess of "tetrahedral" magnetic "moments"). The Me³⁺ ions are magnetized by the weak field of sublattices d and a, and in consequence the moments of the c sites are directed opposite to the magnetization of sublattice d. The distribution of sublattice magnetizations can be pictured schematically as follows:

$$\xrightarrow{c} \stackrel{d}{\xrightarrow{\operatorname{Me}^{3+}}} \xrightarrow{c} \stackrel{d}{\xrightarrow{\operatorname{Fe}^{3+}}} \xrightarrow{a}$$

This scheme provides a possibility of calculating the magnetic moments (per mole) of the garnet ferrites at 0°K. Thus for the case of gadolinium garnet, $3Gd_2O_3 \cdot 5Fe_2O_3$, we have (noticing that $M_c = 7\mu_B$ and $M_d = M_a = 5\mu_B$):

$$6M_c - (6M_d - 4M_a) = 6 \times 7\mu_B - (6 \times 5\mu_B - 4 \times 5\mu_B) = 32\mu_B.$$

Experiment gives $30\mu_{\rm B}$. In the case of yttrium garnet, since Y is nonmagnetic, we have $6M_{\rm d} - 4M_{\rm a} = 10\mu_{\rm B}$, which also agrees with experiment. Thus analysis of the magnitudes of the magnetic moments in garnet ferrites corroborates Néel's "three-sublattice" hypothesis. Recently this hypothesis has been corroborated by neutron diffraction experiments.⁸

Thus in the case of the garnet ferrites we encounter a new form of uncompensated antiferromagnetism, caused by the nonequivalence of three magnetic sublattices. At the present time the following data are known on the magnetic properties of garnet ferrites. The temperature dependence of the spontaneous magnetization has been studied in the range 2.2 to 750°K, by the method of removal from the field, for the following garnet ferrites: $3Me_2O_3 \cdot 5Fe_2O_3$, where Me = Gd, Tb, Dy, Ho, Er, Tu, Yb, Lu, Sm, Eu, or Y. For the majority of the ferrites, the dependence of magnetization on field satisfies the following relation:

 $\sigma_H = \sigma_s + \chi H$,

where $\sigma_{\rm S}$ is the spontaneous magnetization, and χ is the susceptibility in strong fields and is independent of the field. The spontaneous magnetization $\sigma_{\rm S}$ was found by extrapolation to zero field of the straight part of the curve in strong fields.



FIG. 1. Temperature dependence of the spontaneous magnetization $\sigma_{\rm s}$ (in $\mu_{\rm B}$ per mol) for garnet ferrites.

In Figs. 1 and 2 are shown curves of the temperature dependence of the spontaneous magnetization (in μ_B per mol) for all the garnet ferrites investigated. In these figures there are two characteristic peculiarities that should be noted:

1. In the majority of the garnet ferrites (with the exception of the ferrites of Y, Lu, Sm, and Eu), there are observed compensation points, at which the magnetizations of the sublattices compensate each other.

2. It is a surprising and interesting fact that all these substances have about the same Curie point (in contrast to this, the molar magnetic moments at 0°K vary over a relatively wide range). The coincidence of the Curie points must be attributed to



FIG. 2. Temperature dependence of the spontaneous magnetization $\sigma_{\rm s}$ (in $\mu_{\rm B}$ per mol) for garnet ferrites of samarium and europium.

FIG. 3. Dependence of the spontaneous magnetization σ_0 (in $\mu_{\rm B}$ per mol) on content of substitutional Al^{3+} ions.



the fact that the Me^{3+} ions take little part in the ferromagnetism; the interaction of the c lattice with the d and a is small, therefore here the magnitude of the Curie point is determined essentially by the strength of the interaction between the iron ions (the number of which remains approximately the same in all the garnet ferrites).

In the foreign literature there have recently⁹⁻¹¹ appeared numerous papers on the investigation of

FIG. 4. Dependence of the spontaneous magnetization σ_0 (in μ_B per mol) on content of substitutional Cr^{3+} and Ga^{3+} ions.



the replacement of iron ions by other ions in the stoichiometric garnet ferrite $3Me_2O_3 \cdot 5Fe_2O_3$.

Substitutions of this type are directed at two goals: (1) investigation of garnet ferrites with new magnetic properties; (2) elucidation of the role of the sublattices in ferromagnetism; this can be accomplished by changing the number of magnetic ions in the lattices by introducing nonmagnetic or magnetic ions into them.

It should be noted also that experiments of this type are of interest for crystal chemistry. If ions with a different electron configuration and atomic radius are introduced into the sublattices, these Magnetic moment $n\mu_B$ and Curie temperature θ for $3Y_2O_3 \cdot xMe_2O_3 (5-x) Fe_2O_3$

x Me ₂ O ₃	nμ _B observed	θ , ^o K observed
$\begin{array}{c} 0\\ 0.25\ {\rm Ga_2O_3}\\ 0.75\ {\rm Ga_2O_3}\\ 0.34\ {\rm Al_2O_3}\\ 1.00\ {\rm Al_2O_3}\\ 0.25\ {\rm Sc_2O_3}\\ 0.75\ {\rm Sc_2O_3}\\ 0.50\ {\rm In_2O_3}\\ 0.25\ {\rm Cr_2O_3}\\ 0.25\ {\rm Cr_2O_3}\end{array}$	$\begin{array}{r} 9.92 \\ 7.95 \\ 4.35 \\ 7.00 \\ 3.25 \\ 11.97 \\ 14.4 \\ 13.8 \\ 10.95 \end{array}$	545 519 460 497 415 500 365 444 515

ions will prefer fourfold, sixfold, or eightfold sites. and consequently it is possible thus to change the distribution of atoms in the garnet structure and to study their role in magnetization. Most of the experiments have been done by replacement of the iron in yttrium garnet ferrite by the ions: Al^{3+} , Ga^{3+} , Cr^{3+} , Sc^{3+} , In^{3+} . Introduction of Al^{3+} and Ga³⁺ ions decreases the magnetic moments of the garnet ferrites (cf. Figs. 3 and 4 and the table). This must be explained by supposing that Al^{3+} and Ga³⁺ are distributed on tetragonal sites (in sublattice d), in consequence of which the resultant magnetic moment decreases. Al^{3+} and Ga^{3+} have smaller ionic radii than Fe³⁺, therefore they have an affinity for the tetrahedral positions. Sc^{3+} and In^{3+} , having larger radii than Fe³⁺, preferentially occupy octahedral positions, therefore the magnetic moment of the ferrites increases (cf. table). The Cr^{3+} ion, though also smaller than the Fe³⁺, nevertheless occupies octahedral positions, perhaps by virtue of its peculiar electronic configuration. This leads, in a certain Cr^{3+} concentration range, to an increase of magnetic moment.

Thus it follows from the experiments on substitution that the moment of a garnet ferrite can be increased by replacing an iron ion in sublattice a by any nonmagnetic ion, and can be decreased by replacing an iron ion in sublattice d by a nonmagnetic ion. It should be noted that in all cases of replacement of an iron ion by another ion, the magnitude of the Curie temperature of the garnet ferrite decreases. This is explained as a consequence of the fact that the number of magnetic $Fe^{3+}-O^{2-}-Fe^{3+}$ interactions on a magnetic ion in the garnet ferrite decreases. Summarizing everything that has been said about the magnetic properties of garnet ferrites, we must point out that we still have only scanty data on the magnetic properties of these interesting materials. Up to now nothing is known about hysteretic, magnetostrictive, electrical, and other properties of the garnet ferrites. Meanwhile these data are necessary for an understanding of the nature of their spontaneous magnetization.

⁴ R. Pauthenet, Compt. rend. **242**, 1859 (1956); Aléonard, Barbier, and Pauthenet, Compt. rend. **242**, 2531 (1956); R. Pauthenet, Compt. rend. **243**, 1499 and 1737 (1956); Néel, Bertaut, Forrat, and Pauthenet, Izv. Akad. Nauk SSSR, Ser. Fiz., **21**, 904 (1957) [Columbia Tech. Transl. **21**, 903 (1957)].

⁵S. Geller and M. A. Gilleo, Acta Cryst. 10, 239 (1957).

⁶J. W. Nielsen, J. Appl. Phys. **29**, 390 (1958).

⁷ L. Néel, Compt. rend. 239, 8 (1954).

⁸Bertaut, Forrat, Herpin, and Mériel, J. phys. et radium 18, 32s (1957).

⁹G. Villers and J. Loriers, Compt. rend. 245, 2033 (1957).

¹⁰ R. Pauthenet, J. Appl. Phys. 29, 253 (1958).

¹¹ M. A. Gilleo and S. Geller, J. Appl. Phys. 29, 380 (1958).

Translated by W. F. Brown, Jr.

¹J. F. Dillon, Phys. Rev. **105**, 759 (1957); Calhoun, Overmeyer, and Smith, Phys. Rev. **107**, 993 (1957).

² J. Paulevé, Compt. rend. **244**, 1908 (1957); **245**, 1604 (1957).

³ F. Bertaut and F. Forrat, Compt. rend. 242, 382 (1956).